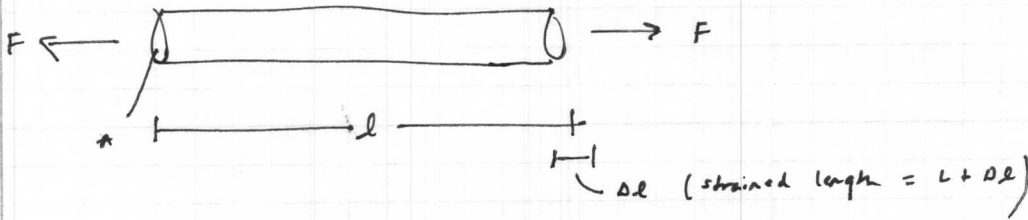


in a macroscopic solid, we have simple linear-elastic behavior

e.g. a bar in uniaxial tension



← young's modulus [N/m^2]

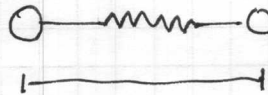
$$\sigma = \frac{F}{A} = E \epsilon$$

$$\epsilon = \frac{\Delta L}{L}$$

$$\frac{F}{A} = E \frac{\Delta L}{L} \Rightarrow \Delta L = \frac{FL}{AE} \quad \text{--- we'll return here later.}$$

at a far smaller scale, let's consider a material as a massive array of springs

→ each bond is like a spring



this is modeled by an interatomic potential,

$x_0 = r_e$
← equilibrium separation

e.g., the Morse potential

$$V = D \left[\left(1 - e^{-\beta(r-r_0)} \right)^2 - 1 \right]$$

↑ well "depth" β = well width.

$$V(x) = \text{energy [J]}$$

$$\frac{dV}{dx} = \text{force [N]}$$

$$\frac{d^2V}{dx^2} = \text{stiffness [N/m]}$$

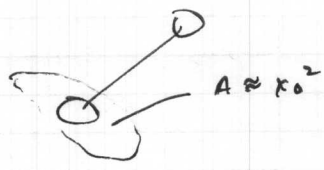
thus, if our "material" is a single bond spring

the restoring force $F = s(x - x_0)$

↑
stiffness.

↘
equilibrium length.

define "area" = $\frac{1}{x_0^2}$



strain = $\frac{\Delta l}{l} = \epsilon$

then stress = $\frac{F}{A} = \frac{F}{x_0^2} = \frac{s(x - x_0)}{x_0^2} = \frac{s}{x_0} \left(\frac{x - x_0}{x_0} \right)$

↖
modulus $[N/m^2] = E$

let's approximate some real values.

based on interatomic potentials.

- covalent bonds, $s = 20 - 200 \text{ N/m}$
 - metal/ionic bonds, $s = 15 - 100 \text{ N/m}$
 - polymers, $s = 0.5 - 1 \text{ N/m}$
- } c-k bonds dominate springs in series

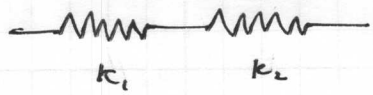
thus, if $E \approx \frac{s}{x_0}$

(c-k bonds are weak)

$x_0 = 0.2 \text{ nm}$

$s = 200 \text{ N/m}$

$E \approx 1000 \text{ GPa} = \text{graphite!}$



$k_{eff} = \left(\frac{k_1 + k_2}{k_1 k_2} \right)^{-1}$

if $k_1 \ll k_2 \Rightarrow k_{eff} = k_1$

stress can typically be close to the theoretical limit

strength, no... it is governed by defects, slip (e.g. grain boundaries)

and can be as low as $10^{-4}E$; in nanoscopic materials, we can reach theoretical strength limits.

strength, lets say bonds break @ 10% elongation.

$$\epsilon = \frac{\sigma}{k_0} \left(\frac{k - k_0}{k_0} \right) \quad \Rightarrow \quad \epsilon \approx 0.1 E$$

$\underbrace{\quad\quad\quad}_{= 0.1}$
 \searrow see Ashby plot.

in polymer fibers, chain alignment is very high \Rightarrow we approach theoretical strength, but are away from strength limit, which is practically governed by chain-chain gaps.

defects are stress concentrators.

